

# Study of Inertial Confinement Fusion via the Equation of State

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- 1. Introduction**
- 2. Why compression**
- 3. Energy gain**
- 4. Bremsstrahlung, degeneracy and the “Clean Fusion” problem**
- 5. Conclusions**

## Ideal gas in a gravitational field

$$\text{internal energy} = E = \frac{PV}{\gamma - 1} = \frac{Nk_B T}{\gamma - 1}; \quad \gamma = \frac{C_P}{C_V}$$

$$\text{kinetic energy} = E_{\text{kin}} = (3/2)Nk_B T = (3/2)(\gamma - 1)E$$

$$\text{virial theorem} \Rightarrow -\sum_{i,j} G \frac{m_i m_j}{|\vec{r}_i - \vec{r}_j|} = E_{\text{pot}} = -2E_{\text{kin}} = 3(1 - \gamma)E$$

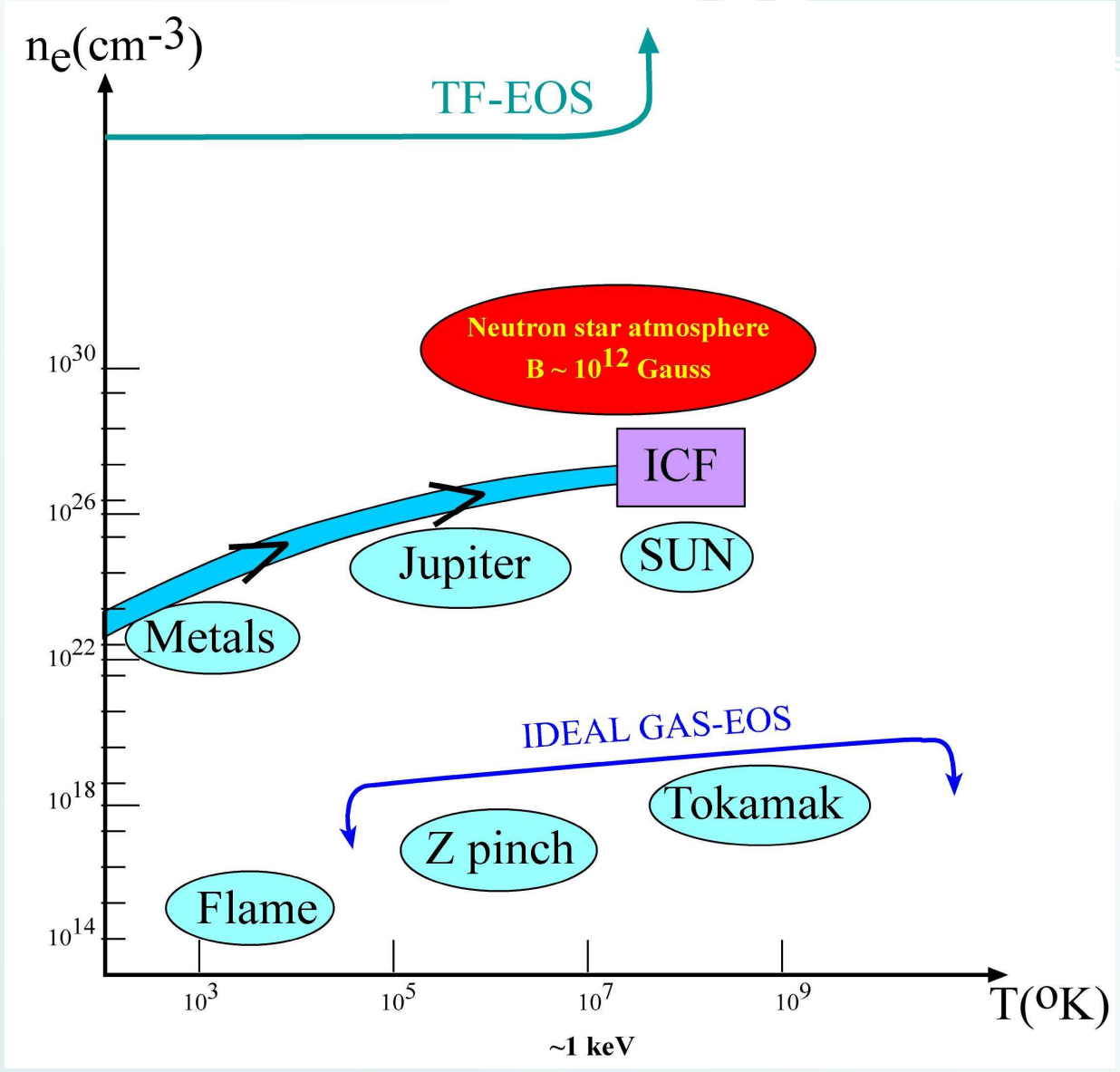
$$E_{\text{total}} = E + E_{\text{pot}} = \frac{(\gamma - 4/3)}{(\gamma - 1)} E_{\text{pot}}$$

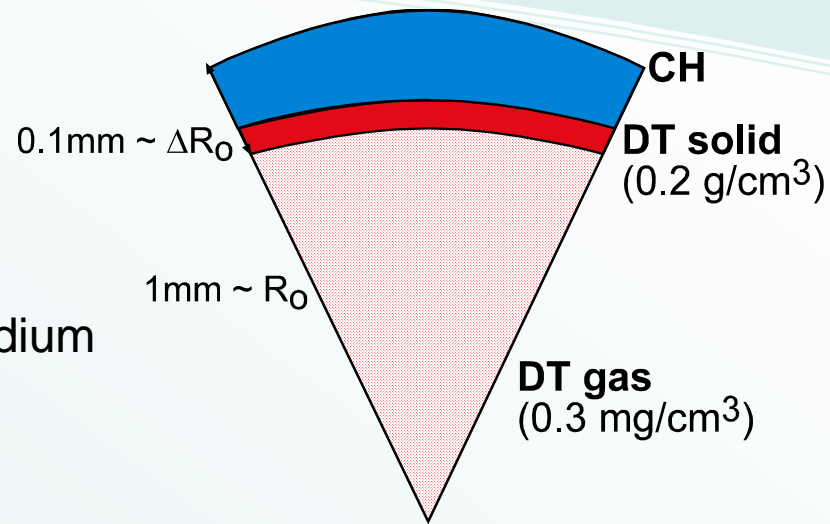
$$E_{\text{total}} = \frac{(\gamma - 4/3)}{(\gamma - 1)} E_{\text{pot}}$$

$\gamma > 4/3 \Rightarrow E_{\text{total}} < 0 \Rightarrow \text{star stable}$

virial theorem  $\Rightarrow G \frac{M}{R} \sim 2E_{\text{kin}} \sim T$

- As the star contracts its internal energy increases and it gets hotter. At the same time it radiates energy!
- This is due to the fact that loss in potential energy is greater than the increase in internal energy and the difference is radiated away.
- For example, for  $\gamma = 5/3$ : only half of the decrease in the potential energy is used in increasing its internal energy (i.e. temperature)

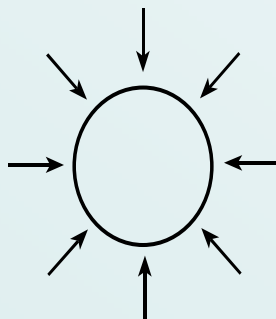




LASER SYSTEM  $\sim$  football stadium

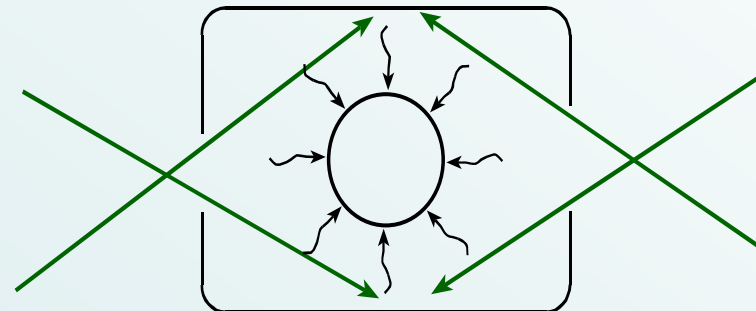
## DIRECT DRIVE

Laser beams shine on a pellet directly



## INDIRECT DRIVE

Laser beams shine on a miniature oven(=hohlraum)  
which  
heats up the pellet by x-rays



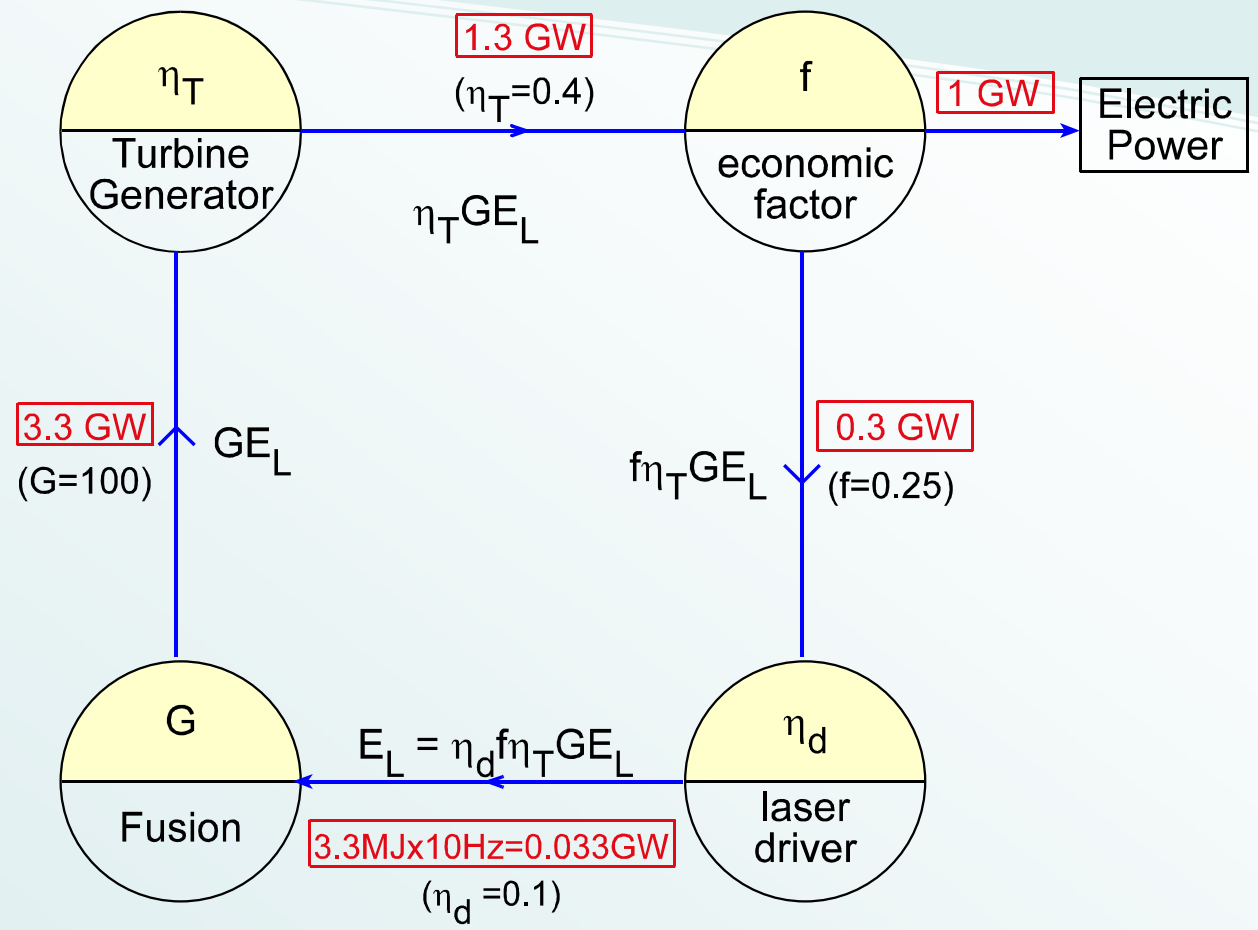
**x-t is space-time,**  
 **$\rho$  = density,  $u$  = flow velocity,**  
 **$P$  = pressure,  $E$  = internal energy**  
**3 hydrodynamic equations,**  
**4 variables ( $\rho$ ,  $u$ ,  $P$ ,  $E$ ),**  
**The fourth equation is EOS!**

$$\text{Mass : } \quad \frac{\partial \rho}{\partial t} = -\frac{\partial}{\partial x}(\rho u)$$

$$\text{Momentum } \quad \frac{\partial}{\partial t}(\rho u) = -\frac{\partial}{\partial x}(P + \rho u^2)$$

$$\text{Energy : } \quad \frac{\partial}{\partial t}\left(\rho E + \frac{1}{2}\rho u^2\right) = -\frac{\partial}{\partial x}\left(\rho E u + P u + \frac{1}{2}\rho u^3\right)$$

$$\text{EOS : } \quad \mathbf{E = E(P, \rho)}$$



$\Rightarrow$   $G \eta_d \eta_T f = 1$

$G \geq 100$

Example:  $\eta_d = 0.1, \eta_T = 0.4, f = 0.25 \Rightarrow G = 100$

## 2. WHY COMPRESSION?

$\eta_H$  = (hydrodynamic) efficiency from laser to plasma

$\eta_A$  = laser absorption by pellet

$$\eta_H \eta_A E_L = (3n_C k_B T) \left( \frac{4}{3} \pi R_C^3 \right)$$

$$E_L = \left( \frac{1}{\eta_H \eta_A} \right) \left( \frac{4\pi}{2.5 m_p \rho_0} \right) \left( \frac{n_0}{n_C} \right)^2 (\rho_C R_C)^3 k_B T$$

using:

$$\rho_0 = 0.2 [\text{g/cm}^3], \quad \rho_C R_C = 3 [\text{g/cm}^2] \quad \{\text{for } \phi \approx 30\%\}$$

$$k_B T = 4 \text{keV} \quad \{\text{for } E(\text{fusion}) > E(\text{bremsstrahlung})\}$$

$$\Rightarrow E_L [\text{MJ}] \approx \left( \frac{1.3 \cdot 10^6}{\eta_H \eta_A} \right) \left( \frac{\rho_0}{\rho_C} \right)^2$$

example:  $\eta_H = 0.1$ ,  $\eta_A = 0.8$ ,

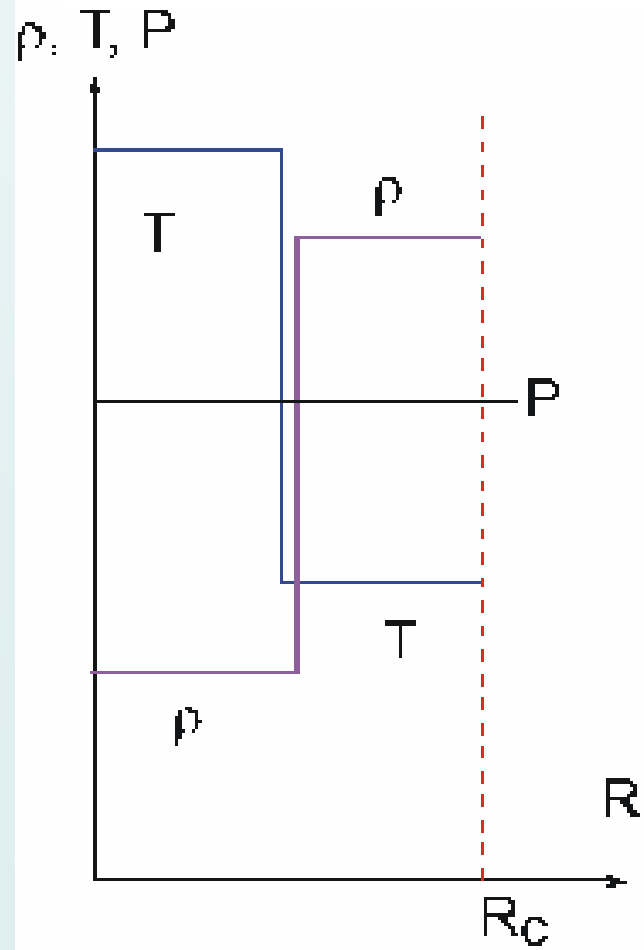
$$\Rightarrow E_L [\text{MJ}] \approx 1.6 \cdot 10^7 \left( \frac{\rho_0}{\rho_C} \right)^2$$

$$\frac{\rho_C}{\rho_0} = 3000 \left\{ \text{i.e. } \rho_C = 600 [\text{g} / \text{cm}^3] \right\}$$

$$\Rightarrow E_L [\text{MJ}] \approx 1.75$$

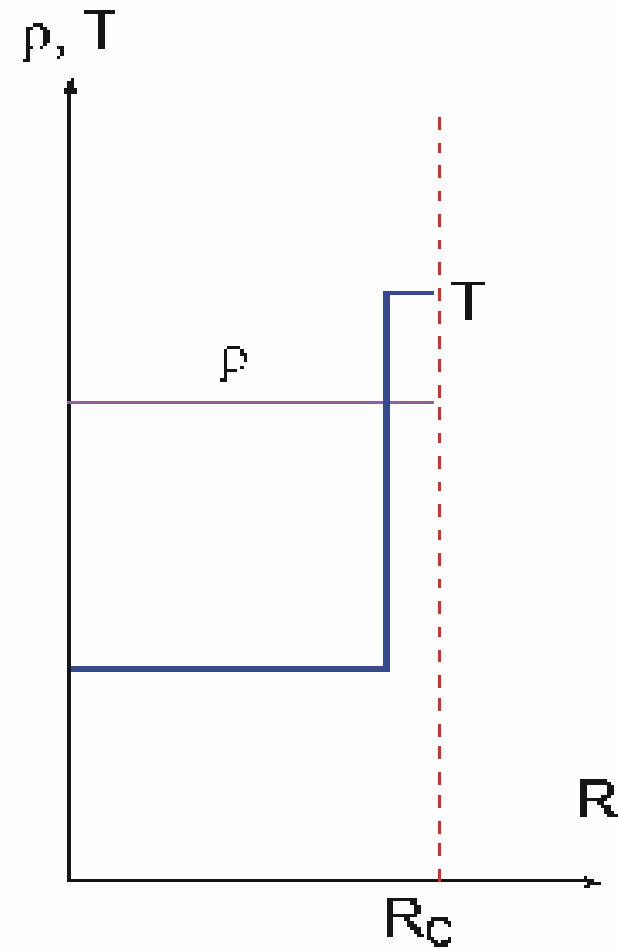
## Spark Ignition

(isobaric:  $P = \text{const}$ )

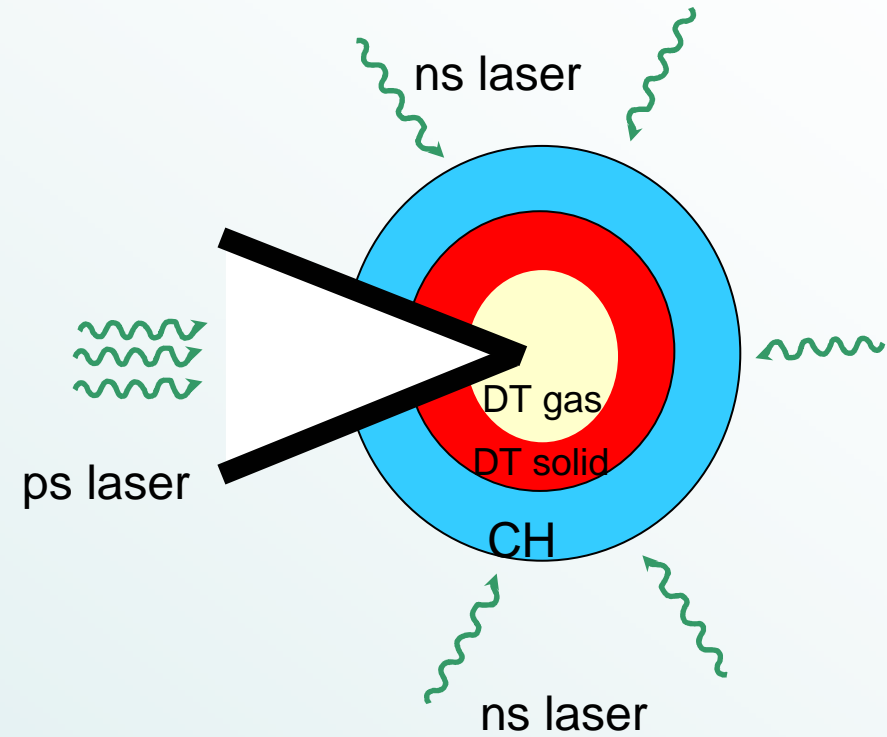
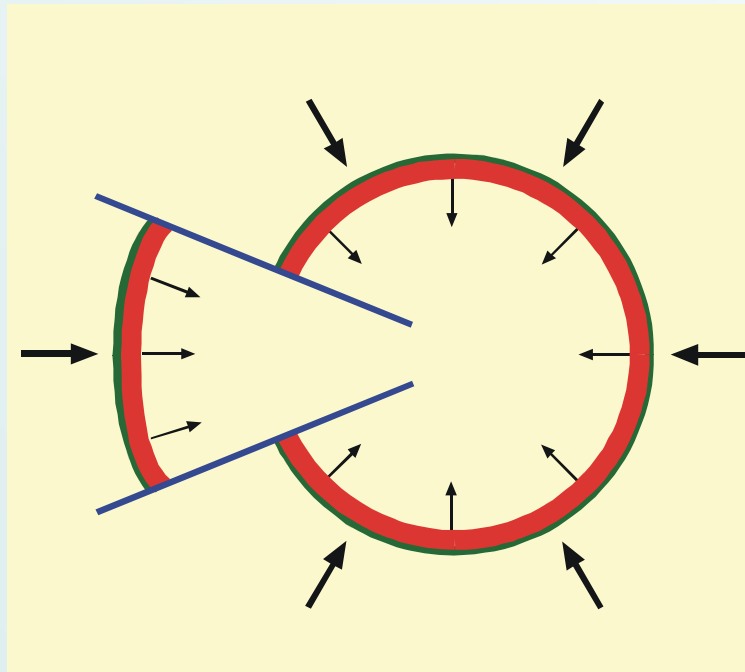


## Fast Ignition

(isochoric:  $\rho = \text{const}$ )



# 3. ENERGY GAIN FAST IGNITION: Impact and Laser



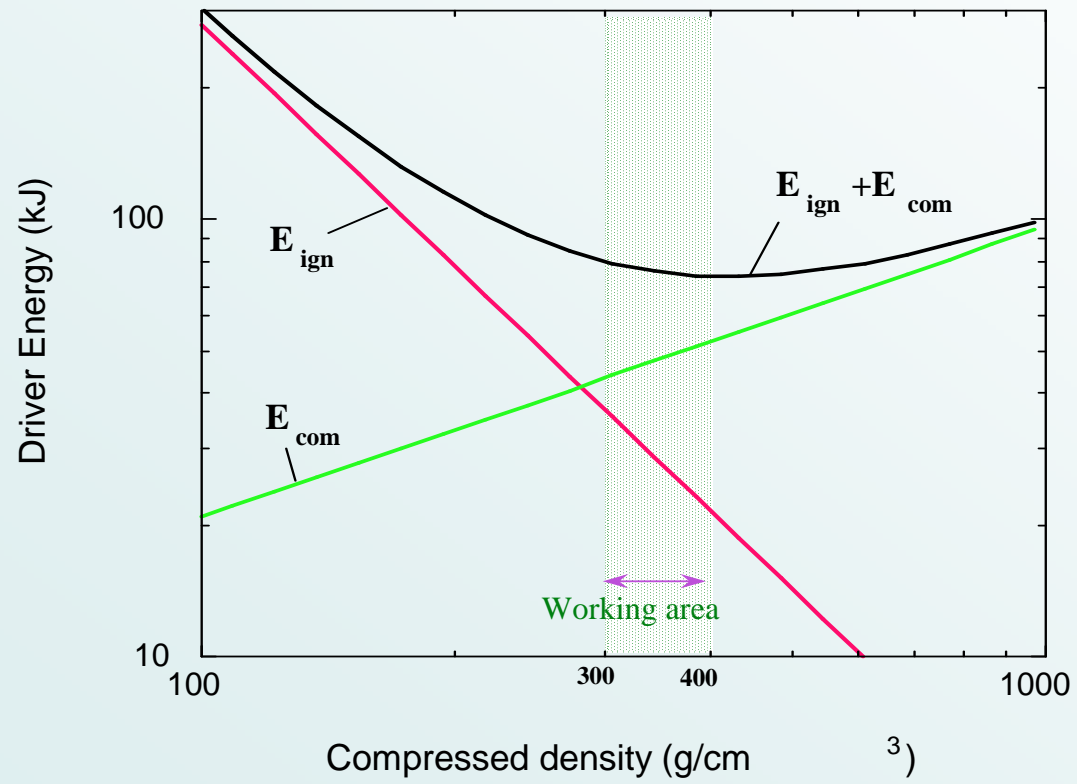
## Gain in Impact Fast Ignition

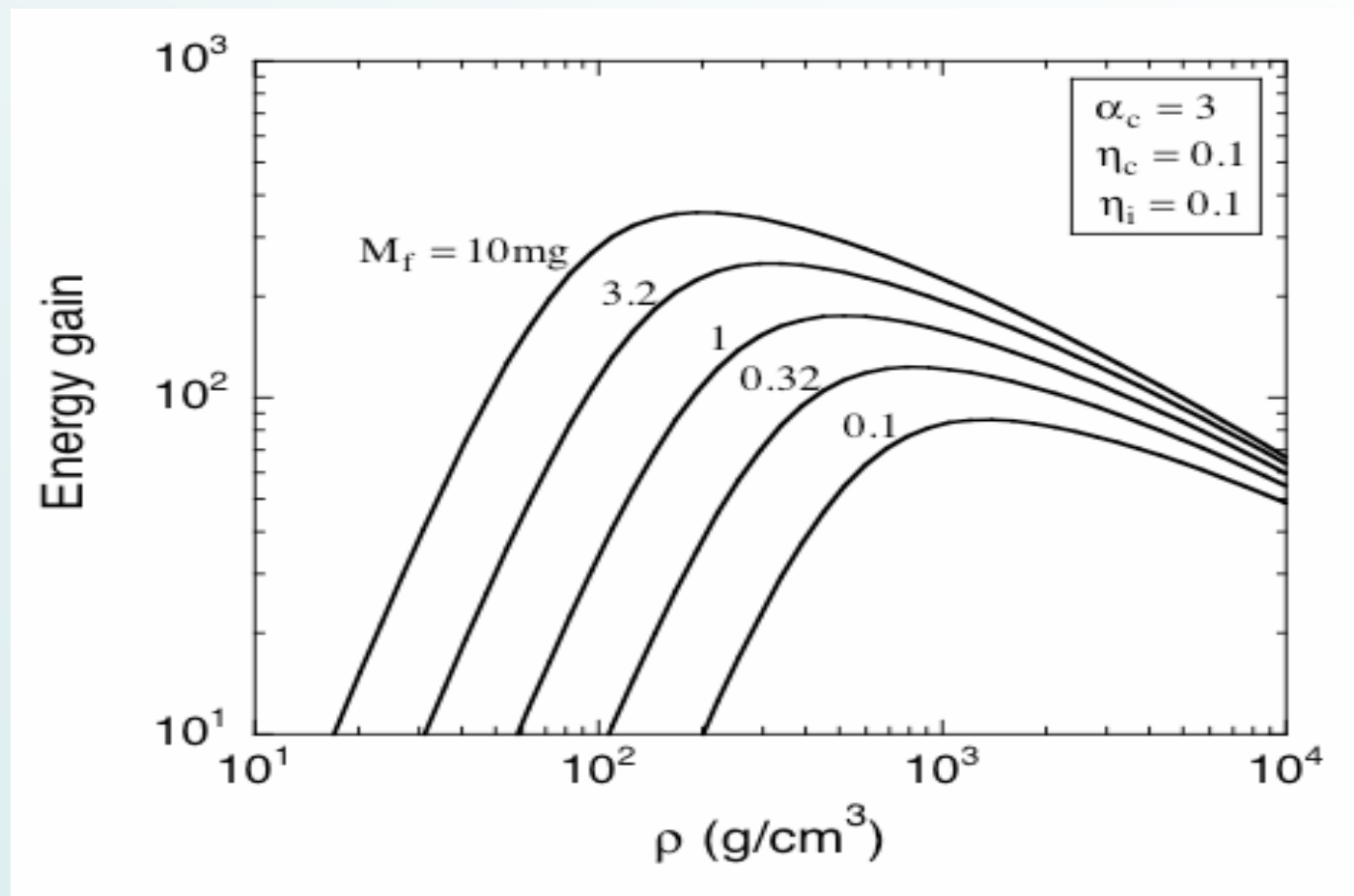
$$E_{\text{ign}} \approx 15[\text{kJ}] \left( \frac{1}{\eta_s} \right) \left( \frac{100\text{g/cm}^3}{\rho_s} \right)^2$$

$$E_{\text{com}} \approx 110[\text{kJ}] \left( \frac{1}{\eta_c} \right) \alpha_c \left( \frac{\rho_c}{\rho_0} \right)^{2/3} \left( \frac{M_c}{\text{g}} \right)$$

$$\text{Total Driver Energy} \equiv E_d = E_{\text{com}} + E_{\text{ign}}$$

$$\text{GAIN} \equiv G = \frac{E_{\text{fusion}}}{E_d} \approx \frac{\phi q_{\text{DT}} M_c}{E_d}$$





## 4. Bremsstrahlung, degeneracy and the $p^{11}\text{B}$ problem

$$\left. \frac{W_{\text{deg}}}{W_{\text{class}}} \right\}_{\text{bremss}} = \left( \frac{\sqrt{\pi}}{2F_{1/2}(\eta)} \right) \left\{ F_1(\eta) - \frac{1}{2} \left[ \ln(e^\eta + 1) \right]^2 \right\}$$

$$F_k = \int_0^\infty \frac{x^k dx}{e^{x-\eta} + 1}$$

$$\eta \equiv \frac{T}{\varepsilon_F} \ll 1 \Rightarrow \left. \frac{W_{\text{deg}}}{W_{\text{class}}} \right\}_{\text{bremss}} = \frac{\pi^2 \sqrt{\pi}}{8} \left( \frac{T}{\varepsilon_F} \right)^{3/2}$$

For the DT case we need a density  $\sim 5 \times 10^4 \text{ g/cm}^3$

for  $p^{11}\text{B}$  the density has to be as large as  $\sim 2 \times 10^7 \text{ g/cm}^3$

$$\epsilon_F = \left(3\pi^2\right)^{2/3} \left(\frac{\hbar^2}{2m_e}\right) n_e^{2/3} \Rightarrow \epsilon_F = \epsilon_{F0} \left(\frac{Z}{Z_0}\right)^{2/3} \left(\frac{\rho}{\rho_0}\right)^{2/3}$$

degeneracy case:  $\epsilon_F \gg T \Rightarrow \epsilon_F \approx 5T$

for DT:  $T \sim 2\text{keV}$

for  $p^{11}\text{B}$ :  $T \sim 100\text{keV}$



$$(a) G_{\max}(\text{DT2e})_{\text{deg}} = \frac{17,600\text{keV}}{2\varepsilon_{\text{F}} + 2 \times (3T/2)} \approx \frac{17,600\text{keV}}{2.6\varepsilon_{\text{F}}} \approx 677$$

$$(b) G_{\max}(\text{DT2e})_{\text{clas}} = \frac{17,600\text{keV}}{4 \times (3T/2)} \approx \frac{17,600\text{keV}}{6 \times 4\text{keV}} \approx 733$$

# Theoretical Maximum Gain for



$$(a) \quad G_{\max} (pB^{11}5e)_{\text{deg}} = \frac{8,700\text{keV}}{5\varepsilon_F + 2 \times (3T/2)} \approx \frac{8,700\text{keV}}{5.6\varepsilon_F} \approx \frac{8700}{2800} \approx 3.1$$

$$(b) \quad G_{\max} (pB^{11}5e)_{\text{clas}} = \frac{8,700\text{keV}}{7 \times (3T/2)} \approx \frac{8,700\text{keV}}{10.5 \times 100\text{keV}} \approx 8.2$$

$$(c) \quad G_{\max} (pB^{11}5e)_{\text{ion only}} = \frac{8,700\text{keV}}{2 \times (3T/2)} \approx \frac{8,700\text{keV}}{300\text{keV}} \approx 30$$

# CONCLUSION 1

**NO CLEAN FUSION!!**

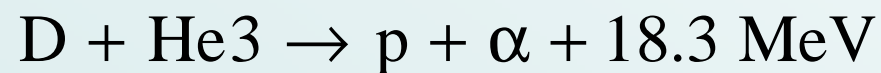
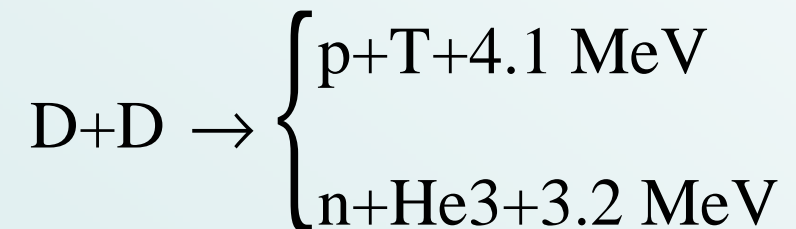


## CONCLUSION 2:

If DT fusion successful (economically)

then we have to master **tritium** and **neutrons**,

even for DD fusion or DHe3 fusion



# Appendix

## SPARK IGNITION

$\xi$  = the fraction of the heated fuel  $\approx 0.14$   
the rest is heated by  $\alpha$  particles  
For  $\rho R \geq 0.3$  [g/cm<sup>2</sup>]  $\alpha$  is absorbed!

$$\Rightarrow G = \frac{\eta_H \eta_A \phi q_{DT}}{\xi q_T} \approx 100$$

PROBLEMS: INSTABILITIES (R-T, etc), MIXING  
 $\Rightarrow$  NO GAIN!

**A SOLUTION: FAST IGNITION**

# THE HOT SPOT

$T_S \approx 5 \text{ keV}$  (fusion energy > bremsstrahlung losses)

$H_S \equiv \rho_S R_S \approx 0.4 \text{ [g/cm}^2\text{]}$  ( $\alpha$  is absorbed for  $0.3 \text{ [g/cm}^2\text{]}$ )

$E_{\text{ign}}$  = laser energy directed into "hot spot"

$\eta_S$  = efficiency from laser to thermal energy

$v_i$  = implosion velocity

$m_i = 2.5m_p$  (i.e. DT fusion)

$M_s$  = total mass of ignitor

kinetic energy = thermal energy

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$M_s$  = total mass of ignitor

kinetic energy = thermal energy

$$\Rightarrow \frac{1}{2} m_i v_i^2 = \frac{3}{2} (k_B T_S) \times 2$$

$$\Rightarrow v_i \approx 1.1 \cdot 10^8 \text{ [cm/s]}$$

$$M_S = \frac{4}{3} \pi R_S^3 \rho_S = \frac{4\pi H_S^3}{3\rho_S^2} = 2.7 \cdot 10^{-5} [\text{g}] \left( \frac{100 \text{g/cm}^2}{\rho_S} \right)^2$$

$$\text{Energy conservation: } \eta_S E_{\text{ign}} = 3 (k_B T_S) \left( \frac{T_S}{m_i} \right)$$

$$\Rightarrow E_{\text{ign}} = \frac{8\pi H_S^3 (k_B T_S)}{5m_p \eta_S \rho_S^2} = 15 [\text{kJ}] \left( \frac{1}{\eta_S} \right) \left( \frac{100 \text{g/cm}^2}{\rho_S} \right)^2$$

# THE MAIN FUEL

$$P_C \equiv \text{fuel (DT) pressure} = n_i T_i + n_e T_e = 2n_e T_C \quad (T_i = T_e \equiv T_C)$$

$n_e$  = electron density,

$$\text{main fuel compressed density} \equiv \rho_C = (2.5 \cdot m_p) n_e$$

$$P_{\text{deg}} \equiv \text{electron degenerate pressure} = \frac{2}{5} n_e \epsilon_F$$

$$\epsilon_F = \text{Fermi Energy} = \frac{(3\pi^2)^{2/3} h^2 n_e^{2/3}}{8\pi^2 m_e} \approx 2.25 \cdot 10^{-11} \rho_C^{2/3} \text{ [cgs, for DT]}$$

$$\alpha_C \equiv \frac{P_C}{P_{\text{deg}}} = \frac{5T_C}{\epsilon_F} \text{ [} \approx 3 \text{ is a reasonable number]}$$

[ $\alpha_C$  is expected to be as small as possible (isentropic compression) ]

$E_{\text{com}}$  = laser energy directed into the main fuel

$\eta_C$  = efficiency from laser to thermal energy

$$M_C = \frac{4\pi\rho_C R_C^3}{3} = \frac{4\pi H_C^3}{3\rho_C^2}, \quad H_C \equiv \rho_C R_C$$

ENERGY CONSERVATION:

$$\eta_C E_{\text{com}} = 3 \left( \frac{M_C}{2.5 \cdot m_p} \right) (k_B T_C) \approx 3.3 \cdot 10^{12} \alpha_C \rho_C^{2/3} M_C$$

# GAIN

$$\text{Total Driver Energy} \equiv E_d = E_{\text{com}} + E_{\text{ign}}$$

$$\text{The burn fraction} \equiv \phi \approx \frac{H_C}{H_C + H_0}$$

$$H_C = \left( \frac{3\rho_C^2 M_C}{4\pi} \right)^{1/3}, \quad H_0 = 7 \text{ [g/cm}^2\text{]}$$

burn fraction from hot spot is neglected ( $H_s \approx 0.4 \text{ [g/cm}^2\text{]}$ ;  $\phi_s \approx 0.05$  )

$$q_{\text{DT}} = (17.6 \text{ MeV}) \cdot \left( \frac{6.02 \cdot 10^{23}}{5 \text{ [g]}} \right) = 3.39 \cdot 10^{11} \text{ [J/g]}$$

$$\text{GAIN} \equiv G = \frac{E_{\text{fusion}}}{E_d} = \frac{\phi q_{\text{DT}} M_C}{E_d} + \frac{\phi_s q_{\text{DT}} M_s}{E_d} \approx \frac{\phi q_{\text{DT}} M_C}{E_d}$$

$$\Rightarrow G = G(\alpha_c, \rho_c, \rho_s, \eta_c, \eta_s, E_d); \quad T_s = 5[\text{keV}], H_s = 0.4[\text{g}/\text{cm}^2]$$

$$G = \frac{10^6 \phi \eta_c}{\alpha_c \rho_c^{2/3}} \left( 1 - \frac{1.5 \cdot 10^{15}}{\eta_s \rho_s^2 E_d} \right) \quad [\text{cgs}]$$

$$M_c[\text{g}] = \frac{\eta_c E_d}{3.3 \cdot 10^{12} \alpha_c \rho_c^{2/3}} \left( 1 - \frac{1.5 \cdot 10^{15}}{\eta_s \rho_s^2 E_d} \right)$$

example:  $\alpha_c = 3, \eta_c = \eta_s = 0.1, \rho_c = 300[\text{g}/\text{cm}^3], \rho_s = 200[\text{g}/\text{cm}^3]$

$$E_d = 100[\text{kJ}] = 10^{12}[\text{erg}]$$

$$\Rightarrow M_c = 4.6 \cdot 10^{-4}[\text{g}], \quad G \approx 120$$